

# The microwave properties of simulated melting precipitation particles: sensitivity to initial melting

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## Abstract

A simplified approach is presented for assessing the microwave response to the initial melting of realistically-shaped ice particles. This paper is divided into two parts: (1) a description of the Single Particle Melting Model (SPMM): a heuristic melting simulation for ice-phase precipitation particles of any shape or size. SPMM is applied to two simulated aggregate snow particles, simulating melting up to 0.15 melt fraction by mass; and (2) the computation of the single-particle microwave scattering and extinction properties these hydrometeors, using the discrete dipole approximation (via DDSCAT), at the following selected frequencies: 13.4, 35.6, 94.0 GHz for radar applications; and 89, 165.0 and 183.31 GHz for radiometer applications. These selected frequencies are consistent with current microwave remote sensing platforms, such as CloudSat and the Global Precipitation Measurement (GPM) mission. Comparisons with calculations using variable-density spheres indicate significant deviations in scattering and extinction properties throughout the initial range of melting (liquid volume fractions less than 0.15). Integration of the single-particle properties over an exponential particle-size distribution provides additional insight into idealized radar reflectivity and passive microwave brightness temperature sensitivity to variations in size/mass, shape, melt fraction, and particle orientation.

## 1 Introduction

Present methods of passive and active microwave remote sensing of precipitation have a key problem: the uncertainty of the physical and associated radiative properties of ice- and mixed-phase snowflakes. In nature, ice particles manifest themselves in an extraordinarily diverse variety of sizes, shapes, and habits – ranging from simple crystals such as needles or plates to complex aggregates and rimed particles.

Microwave radiation is sensitive to the presence of liquid water on ice-phase precipitation (e.g., Klassen, 1988), due, primarily, to the large difference in the dielectric constants of the two materials at microwave frequencies. However, the relationship between the early stages

of melting and incident microwave radiation has not been well described in the literature. With this in mind, the present research seeks to quantify the sensitivity of microwave scattering and extinction cross-sections to the onset of melting (melt fractions less than 0.15), with an emphasis on ultimately improving the forward-model simulation of the physical and radiative properties of realistically-shaped mixed-phase precipitation particles. The onset of melting is generally believed to represent the most rapid changes in the scattering and extinction properties of hydrometeors. Bohren and Battan (1982) made direct measurements and simulation of a normalized radar backscattering cross-section for spongy ice spheres with various liquid water contents. Most notable is that they show a rapid change in the backscattering cross-section between 0 and about 7 percent liquid water volume fraction. Willis and Heymsfield (1989) also noted that the onset of melting represents a critical point where the radar reflectivity increases rapidly. A comprehensive examination of the microwave sensitivity to the entire range of melting is left for future research.

The model described for the first time here, named the Single Particle Melting Model (hereafter SPMM), is a heuristic model designed to provide a basis for simulating the physical description of melting individual ice crystals having an arbitrary shape. By using simple rules and nearest-neighbor interactions, the melting process is simulated with a reasonable facsimile of reality. In SPMM, there are no explicit thermodynamic or physical properties, other than the 3-D shape, and the relative positions of liquid and ice constituents. SPMM is an extremely computationally efficient algorithm for creating a series of melted particles ranging from unmelted to completely melted, only requiring several minutes on a normal desktop computer for a single particle. Thermodynamic melting-layer models can easily be employed to determine bulk melt-water generation (Mitra et al., 1990; Olson et al., 2001), and the SPMM-melted particles can be mapped into an appropriate particle size distribution according to the layer-averaged melting properties.

While the general physical properties and thermodynamics of melting snowflakes is fairly well understood; the complex interaction between realistically-shaped melting snowflake aggregates and incident microwave radiation has been sparsely examined. Recently, radar properties of melting aggregates at 3.0 and 35.6 GHz have been simulated by Botta et al.

(2010). There are no known simulations of realistically-shaped melting aggregate hydrometeors at frequencies above 35.6 GHz, for either passive or active microwave remote sensing applications. For dry realistically shaped aggregates, Petty and Huang (2010) compared simulated reflectivities from aggregates and soft spheres at Ku, Ka, and the passive microwave response 18.7, 36.5, and 89.0 GHz; and Kulie et al. (2014) has examined the Ku, Ka, and W band (14, 35, and 94 GHz) response to variations in shape and size, but neither paper consider melting.

In this study, the discrete dipole approximation (DDA), using DDSCAT version 7.3 (Draine and Flatau, 1994), is employed to compute the scattering and extinction efficiencies of individual particles with mass-melt fractions ranging from from 0.0 (unmelted) to 0.15 (lightly melted). Due to the significant computational requirements of DDSCAT when liquid water is present, we have limited this study to two selected aggregate particles, shown in Fig. 1. The chosen microwave frequencies – 13.4, 35.6, 89.0, 94.0, 165.0, and 183.31 GHz – are relevant to current passive and active microwave sensors, such as CloudSat (Stephens et al., 2002), the recently launched Global Precipitation Measurement mission (GPM, Hou et al., 2014), and instrumentation employed during aircraft- and ground-based precipitation validation experiments over the past decade (e.g., MC3E Petersen and Jensen, 2012, GCPEX Hudak et al., 2012, and IPHEX Barros, 2014) and the upcoming OLYMPEX field campaign in Winter 2015/16 (McMurdie et al., 2015).

The following sections describe the melting simulation methodology, the single-particle scattering and extinction properties at standard radar and radiometer center frequencies, and the particle size distribution averaged properties with implications for remote sensing applications.

## 2 Physical description of melting

In an actual melting snowflake, the distribution of water on the surface of a melting ice-crystal is governed primarily by the local amount of liquid water. According to Oraltay and Hallett (2005), in the initial stages of melting, melt-water distributes more-or-less evenly over

the surface of the ice crystal. At a locally-critical liquid water content, surface tension effects take over and convex water droplets form at branch points and other energetically favorable regions on the surface of the underlying ice-crystal. As the crystal shape changes due to additional melting, these water droplets coagulate and tend to collect toward a common center under the influence of surface tension. Previous studies have examined the effects of modifying the distribution of meltwater on oblate spheroids (Tyynelä et al., 2014), and one has explicitly examined melting realistically shaped snowflakes, but with ad hoc placement of water spheres to simulate melting (Botta et al., 2010). In reality, tumbling, breakup, shedding, collision/aggregation, evaporation/condensation, and other physical interactions often significantly alter growth and melting processes. Due to the complexities of such simulations, we will do not explicitly consider these effects in the present physical model – although aggregation is simulated through the creation of aggregate snowflakes.

## 2.1 Single Particle Melting Model (SPMM)

The Single Particle Melting Model (SPMM), developed by the author and described for the first time here, performs physical particle simulations on an integer-indexed three-dimensional Cartesian grid. Each occupied point represents a small, but finite unit of mass of either ice or liquid. This assemblage of ice and water points constitutes the particle mass and volume; Fig. 1 illustrates this for the two selected aggregate shapes used in this study. In this figure, blue regions represent solid ice and red regions represent meltwater (red is used for visual clarity). The maximum melt fraction is 0.15 for the present study, in order to assess the microwave sensitivity to the initial stages of melting. Throughout this melting range, casual inspection reveals that very little structural changes occur in either the dendrite aggregate (DA) or needle aggregate (NA).

The basis for melting and melt-water movement in SPMM occurs through nearest-neighbor interactions. In a 3-D domain, any given point has 26 nearest neighbors (including all diagonals). The interaction distance is limited to one neighbor, simplifying the computational requirements of the algorithm. Figure 2 conceptually illustrates the melting process in

two dimensions; the same logic applies to the three-dimensional model. The melting simulation proceeds iteratively until all ice is melted and a nearly spherical droplet is formed.

The SPMM proceeds with the following steps:

1. Populate a 3-D Cartesian grid with “ice” points, the ensemble of which comprise the entire volume of the simulated snowflake or aggregate.
2. Iterate over all ice points, tabulating all 26 nearest neighbors (6 sides, 23 diagonals, excluding self).
3. The ice points having the fewest ice neighbors are melted (see Fig. 2a). A stochastic control factor is employed to control the rate of melting.
4. After each melt iteration, a movement check is applied to those liquid points having zero ice neighbors. Prohibiting the movement of liquid points that do have ice neighbors simulates a “coating” effect; whereas liquid points with no ice neighbors are able to move (Fig. 2b and c).
5. Movement is a weighted random walk, subject to certain constraints. The walk is weighted toward total particle center of mass, simulating the coalescence of liquid water. The movement phase iterates until all moving liquid cannot move to an open space closer to the center of mass than the current position. Return to step (2).

In this simple model, ice structure collapse, breakup, or water shedding are not explicitly simulated – any orphaned droplets created during melting will naturally migrate towards the total center of mass as cohesive droplets.

Although the present study considers melt fractions up to 0.15, SPMM provides melting simulations for any arbitrary particle shape until it is completely melted. In principle, it can be applied to the melting of any material, where surface tension is a dominant factor in the liquid phase. The melting increments and distribution of melt-water can be finely-tuned to suit the application. On a modern desktop computer, a particle having 200 000 ice points (“dipoles” in DDA parlance) can complete the entire melting process in less than 5 min using

a single processor core. The Fortran 90 codes and MATLAB codes for SPMM are freely and indefinitely available upon request from the primary author.

### 3 Single-particle scattering and extinction properties

In the previous section we described the method for simulating the melting of realistically shaped precipitation hydrometeors using SPMM. Given these two example sets of melted ice particles, the single-particle extinction cross-section, backscattering cross-section, and asymmetry parameter properties were computed at 23 melting steps between 0.0 (unmelted) to 0.15 (lightly melted, see Fig. 1). To compute the single-particle radiative properties, the DDSCAT 7.3 model is used (Draine and Flatau, 1994).

#### 3.1 Numerical scattering calculations

DDSCAT is an implementation of the discrete dipole approximation (DDA) method for solving Maxwell's equations for linearly-polarized electromagnetic plane waves incident upon an arbitrarily shaped dielectric material consisting of up to 9 different dielectric materials, here we only use two discrete constituent materials: ice and liquid water, neglecting the dielectric constant of air and water vapor. For satisfactory convergence, DDSCAT requires that  $|m|kd < 0.5$ , where  $|m|$  is the magnitude of the complex index of refraction of ice (Warren and Brandt, 2008) or liquid water (Liebe et al., 1991),  $k$  is the wavenumber of the incident radiation, and  $d$  is the minimum spacing between adjacent "dipoles". In each of the following calculations, specific care was taken to ensure that this convergence criteria was adequately satisfied. In practice this was accomplished by having a sufficiently large number of dipoles representing the shapes, minimizing the dipole spacing  $d$  compensating for the increase in  $|m|$  imposed by the increasing amount of liquid water present in the melting simulation.

DDSCAT requires the following inputs: The polarization and wavelength of incident radiation, the 3-D Cartesian position and index of refraction for each "dipole" point, the effective radius of the entire particle (aeff: the radius of a sphere having equal mass), and 3-D ro-

tation angles of the target in the reference frame. Here, the effective radius ( $a_{\text{eff}}$ ) acts as a proxy for particle mass, independent of the particle shape.

The outputs of interest for this study are extinction cross-section, scattering asymmetry parameter, and radar backscattering cross-section at 30 effective radius intervals – ranging in log scale from 50 microns to 2500 microns. For each effective radius value, the input shape (relative dipole positions) remain the same, but each dipole mass (and consequently effective volume) is scaled appropriately. We also examine the extinction and backscattering efficiencies,  $Q_{\text{ext}}$  or  $Q_{\text{bck}}$ ; which are simply the respective cross-sections (e.g.,  $C_{\text{ext}}$ ) divided by the cross-sectional area of the equivalent sphere (defined in DDSCAT by  $Q_* = C_*/\pi a_{\text{eff}}^2$ ).

It should be noted that scaling particles in this manner does not create a mass-dimension relationship that is consistent with observations when these particles are used together in an ensemble. Each particle and its associated scattering and extinction properties are intended to be considered as independent particles if they're being used in an actual retrieval framework, and it's up to the researcher to ensure that the mass-dimension relationship they prefer to use is followed. Leinonen and Moishev (2015) found that the method of producing aggregates can result in significantly different optical properties, depending on the ultimate mass-dimension relationship. In the present study, the purpose of choosing the shape-preserving scaling approach was so that the melting morphology would be preserved for a given base shape across all ranges of masses, enabling a consistent comparison independent of shape changes.

To simulate randomly oriented hydrometeors, an average over multiple orientations of the aggregate relative to a fixed direction of incident radiation is computed. This provides an orientation-averaged set of scattering and extinction properties. Although not shown here, our sensitivity studies suggest that 75 discrete orientations, sampling a full 3-D rotation, is sufficient to provide a reasonably precise orientationally-averaged set of scattering and extinction calculations. This trade-off keeps the computational requirements tractable: for a single effective radius, single shape, and single frequency, it requires 24 h to perform 1 set of calculations running in a parallel implementation a 24-core Intel(R) Xeon(R) CPU X5670 at 2.93 GHz, and requires up to 32 GB of allocated RAM per DDSCAT process. For



6 frequencies, 2 shapes, and 30 effective radii results in computation times on the order of 200 days of continuous calculations (with 3 processes running in parallel). In addition the orientation-averaged quantities, the scattering and extinction properties are also tabulated for each individual orientation, which has important implications for exploring the polarization of scattered radiation, but this is left for future research.

In the current study, the single particle scattering and extinction calculations are divided into two-frequency groups:

1. Radar-specific frequencies, approximately consistent with the Global Precipitation Measurement mission Dual-Wavelength Precipitation Radar (GPM DPR) at Ku-band and Ka-band (13.4 and 35.6 GHz), and with the CloudSat W band radar (94.0 GHz). Several ground-based and aircraft-based instruments also use these channels for precipitation remote sensing applications.
2. “High-frequency” passive microwave channels at 89, 165, and 183.31 GHz, consistent with the GPM Microwave Imager (GMI) channels used for detecting scattering by ice clouds and ice-phase precipitation. Similar channels are used on other satellites, aircraft, and ground-based passive microwave instruments (e.g., SSMI/S, TRMM, CoS-MIR) for cloud and precipitation remote sensing.

### 3.2 Backscattering and extinction at radar frequencies

In Fig. 3, the 13.4 GHz extinction and radar backscattering cross-sections are computed for the dendrite-aggregate (DA) in panels (a) and (b) respectively; and for the needle-aggregate (NA) panels (c) and (d). Colors represent the particle effective radius (values on the colorbar are in microns), which is directly related to the particle mass. Both extinction and backscattering cross-sections are most strongly influenced by changes in size. At the smaller sizes, however, the onset of melting has a strong influence on extinction (panels a and c). This indicates that the onset of melting is characterized by a rapid increase in extinction, while the backscattering tends to exhibit a more linear increase. This is consistent with the integrated backscattering and extinction properties, shown in section 4.

Figure 4 shows the single-particle DDSCAT calculations of the extinction efficiency ( $Q_{\text{ext}}$ ) for the dendrite-aggregate (DA) and needle-aggregate (NA) at 50 micron effective radius for (a) 13.4, (b) 35.6, and (c) 94.0 GHz. Shaded regions represent the range of variations in  $Q_{\text{ext}}$  due to the range of orientations of the particle. Black lines are the equal-mass ice sphere  $Q_{\text{ext}}$  values at 1, 10, 50, and 100 % densities (100 % is equivalent to a density of  $917 \text{ kg m}^{-3}$ ). The extinction increases rapidly with the onset of melting. For every 0.05 increase in melting fraction, the extinction nearly doubles at all three frequencies. Both the needle aggregate and dendrite aggregate extinction efficiency exhibit roughly similar behavior, but is obviously different from the spherical properties, and often outside of the range that could be captured by spherical particle shapes alone (e.g., Liao et al., 2013).

Figure 5 is the same as Fig. 4 except that the effective radius has increased to 2500 microns. The behavior at (a) 13.4 and (b) 35.6 GHz shows a similar sensitivity to the onset of melting, but not quite as rapid as the smaller particle case. Panel (c) 94 GHz shows a relative insensitivity of extinction to onset of melting, in this case it is due to a trade-off between a rapidly decreasing scattering contribution and equally increasing absorption contribution to the total extinction. Also of note in panel (c) is that the needle aggregate now shows an overall larger extinction than the dendrite aggregate, different from the other frequencies. This is indicative of the increased sensitivity of the smaller wavelength (approximately 3.2 mm at 94 GHz) to the finer-scale structures present in the needle aggregate. We believe that this is the first evidence in the existing literature that points towards this behavior, however additional research will be required on this topic.

Following the same approach as Figs. 4 and 5, Figs. 6 and 7 present the backscattering efficiency, which is the differential cross-section for backscattering ( $\mu\text{m}^2 \text{ sr}^{-1}$ ) divided by  $\pi a_{\text{eff}}^2$  (Bohren and Huffman, 1983). Integrating the single-particle backscattering cross-section over a distribution of particle sizes yields the radar reflectivity, which is discussed in Sect. 4. Figure 6 is the backscattering efficiency for particles of a 50 micron effective radius, whereas Fig. 7 is the same at 2500 micron effective radius.

At 50 micron effective radius, the scattering is well into the Rayleigh regime at all of the radar frequencies considered here. The response is a relatively gentle increase in backscat-

tering over the 0 to 0.15 melt fraction range. The 94 GHz backscattering efficiency is roughly 3 orders of magnitude larger than the 13.4 GHz efficiency. Also of note is that the variable density spheres (black lines) consistently underestimates the total radar backscattering for all densities compared to the two dendrites, implying that no modification of the density could reproduce the backscattering obtained by the non-spherical particles.

Figure 7 at 2500 microns shows a more complex relationship to melting and particle orientation. The backscattering efficiency exhibits a large variance due to particle orientation (the shaded regions) in all panels. The spheres at 13.4 GHz encompass the range of variability and exhibit a general increase in backscattering with melting. However this breaks down at 35.6 and 94 GHz, where the Mie resonance effects start to have a stronger influence on the computed backscattering. Similar to what was seen in Fig. 5, the reversal of the backscattering roles occurs at 35.6 and 94 GHz, with the needle aggregate exhibiting a higher backscattering than the dendrite aggregate. It is also notable that at 94 GHz, the backscattering decreases with increasing melt fraction.

### 3.3 Scattering and extinction at passive microwave frequencies

At passive microwave frequencies commonly employed for snowfall retrieval (e.g., 89, 165, and 183.31 GHz), the particle interaction with microwave radiation comes primarily through extinction (scattering + absorption) and emission. Due to space constraints, only one effective radius (1500 microns) is shown here in order to understand the general sensitivity of the single-particle extinction, single scattering albedo and asymmetry parameter to the onset of melting.

In Fig. 8 we have computed the extinction efficiency for the needle aggregate (NA) and dendrite aggregate (DA) vs. melt fraction. Generally the spheres do not adequately cover the observed ranges of extinction for the non-spherical particles. At 89 GHz (panel a), the DA extinction is higher than NA, this role reverses at 165 and 183.31 GHz. Roughly speaking 165 and 183.31 GHz exhibit similar sensitivities to melting at all effective radii (only 1500 micron is shown here); suggesting that for rough estimates, computing the 165 GHz properties may be sufficient for capturing the scattering and extinction behaviors at 183.31 GHz

(and nearby channel offsets employed on the GPM microwave imager, and other imagers and sounders). In actual remote sensing applications, the difference in water vapor emission/absorption at 165 vs. 183.31 GHz is likely to dominate the signal, except in the driest of atmospheric profiles (Skofronick-Jackson and Johnson, 2011).

5 The single scattering albedo (the ratio of scattering to the total extinction) presented in Fig. 9, tells us the primary story of interest with regard to the onset of melting. Specifically we see that a small change in melt fraction yields a significant linear decrease in single scattering albedo, which is an indicator of the rapidly increasing contribution of absorption to the total extinction. This, in turn, drives the thermal emission (according to Kirchhoff's law) that we will later observe in Sect. 4.

10 Finally, in Fig. 10, the scattering asymmetry parameter (cosine averaged scattering contribution over all angles) describes the degree to which incident radiation is forward scattered ( $g > 0$ ) or backward scattered ( $g < 0$ ). At a 1500 micron effective radius, very little sensitivity to the onset of melting was observed. Similar results were found at other effective radii. This is consistent with the notion that the scattered radiation depends primarily on the shape of the particle. Over this range of melting, the actual particle shape has changed very little, so the degree of forward scattering is not expected to change much. If melting were to continue beyond a melt fraction of 0.15 (not shown), the asymmetry parameter would change as the shape of the particle changes.

20 Overall, the single scattering properties show a marked sensitivity to the onset of melting for scattering and extinction, with the exception of the asymmetry parameter. The spherical particle approximation does not produce scattering and extinction properties that have similar behavior to the non-spherical particle, particularly as the frequencies change. In some cases, the spherical particle properties do not bracket the non-spherical particles' properties, suggesting that under the current formulation, no amount of modifying the density parameter could result in a reliable substitute for more physically realistic shapes when one considers all of the scattering and extinction properties of interest to passive and active remote sensing applications.

## 4 Integrated properties

In the previous section, we examined a subset of the single-particle scattering and extinction properties. Of interest for atmosphere radiative transfer and remote sensing applications is how these quantities behave in an ensemble of particles.

### 4.1 Radar Response

The equivalent radar reflectivities ( $Z_e$ ) at each frequency are calculated using the single-particle radar backscattering cross sections,  $C_{\text{bck}}(D)$ , and integrating these over a given particle size distribution,  $N(D)$ :

$$Z_e = \frac{\lambda^4}{\pi^5 |K_w|^2} \int_{D_{\text{min}}}^{D_{\text{max}}} C_{\text{bck}}(D) N(D) dD, \quad (1)$$

where  $D$  is the mass-equivalent diameter (twice the effective radius),  $\lambda$  is the wavelength of incident radiation with the same units as  $D$ , and  $K_w = (m_w^2 - 1)/(m_w^2 + 2)$  is the dielectric factor computed from the refractive index of water,  $m_w$ . We note that in the present case, we specifically chose  $|K_w|^2 \approx 0.93$  for all wavelengths and melt fractions as a comparative convenience, so that only the backscattering properties vary with melt fraction. In actual radar reflectivities, the response to the dielectric properties of the ice / snow / mixed phase particles within the radar range gate will be different, and completely dependent on the individual constituents shape, composition, temperature, and frequency of incident radiation.

For lack of a suitable alternative, we have assumed that the melt fraction of each size of particle is the same fraction across all particles in the distribution. This provides a constant melt-fraction quantity, independent of mass, so that radar sensitivity to variations in the melt fraction can be readily examined without confusion. Future research will explore the variation of melt fraction for particles of different sizes in a given volume of the atmosphere.

For the snowfall particle size distribution, we choose the exponential size distribution from Sekhon and Srivastava, 1970; Johnson, 2007. In each of the following calculations, all

comparisons are made using an equal particle mass distribution. The particle size distribution  $N(D)$  is given by

$$N(D) = N_0 \exp\left(-3.67 \frac{D}{D_0}\right), \quad (2)$$

where  $N_0$  is the intercept parameter, and  $D_0$  is the “characteristic” diameter of the PSD.  $N_0$  and  $D_0$  are related as follows:

$$D_0 = 0.14 R^{0.45} [\text{cm}], \text{ and} \quad (3)$$

$$N_0 = 2.5 \times 10^3 R^{-0.94} [\text{mm}^{-1} \text{m}^{-3}], \quad (4)$$

when liquid-equivalent precipitation rate,  $R$ , is given in units of  $\text{mm hr}^{-1}$ .

Figures 11 and 12 show the simulated equivalent radar reflectivity, averaged over all particle orientations, vs. melt fraction for the needle aggregate (NA) (green line) and dendrite aggregate (DA) (blue line), along with the mass-equivalent variable-density spheres (black lines). In Fig. 11, we have selected the characteristic diameter  $D_0 = 0.11 \text{ cm}$  (corresponding to an ice water content of approximately  $0.13 \text{ g m}^{-3}$ ) for analysis. At all three frequencies, the reflectivity increases by a few dB over the range of melting. With a judicious choice of density, spheres could reasonably approximate the simulated reflectivities of the non-spherical particles, except for the dendrite aggregate at 13.4 GHz (Fig. 11a).

In Fig. 12, the orientation-averaged reflectivities are computed for  $D_0 = 0.55 \text{ cm}$ , corresponding to an ice water content of approximately  $3 \text{ g m}^{-3}$ , which is roughly  $15 \text{ mm h}^{-1}$  liquid-equivalent snowfall rate – an exceptionally high snowfall rate with very large snowflakes, representing an upper limit. Consequently, the simulated equivalent radar reflectivities are as high as 47 dBZ for (DA) at 13.4 GHz. The increase in reflectivity over the onset of melting is, similar to Fig. 11, only a few dBZ from 0 to 0.15 melt fraction.

We note that because  $D_0 = 0.55 \text{ cm}$  is greater than the upper limit of the individual particle size, truncation errors in the particle size distribution integration are potentially significant. So these quantities should be interpreted with this in mind. Due to computational

limitations at the time of research, the individual melt particle sizes cannot be reliably computed above 2500 microns effective radius (5000 microns effective diameter).

Although not shown here, continued melting does not significantly increase the reflectivity beyond this point. This result may appear to be inconsistent with radar observations of melting precipitation (i.e., the radar bright-band) (Klassen, 1988). This suggests that the changing dielectric properties alone do not create the observed rapid increase in radar reflectivities associated with melting. It is postulated, without proof, that enhanced aggregation of particles at or near the melting point causes a rapid increase in total particle size, enhancing the  $D^4$  diameter dependence of the reflectivity factor (e.g., Fabry and Zawadzki, 1995). A more detailed analysis of enhanced aggregation is beyond the scope of the current research.

Figures 13 and 14 are physically consistent with Figs. 11 and 12 respectively, but now show the volume extinction coefficient. The most apparent feature is the rapid increase in extinction with the onset of melting. This behavior was observed in the single particle scattering properties (e.g., Figs. 3–5). The integrated extinction suggests that attenuation of the radar beam starts to accumulate rapidly with only a modest increase in reflectivity. As before, the range of extinction of the spherical particle approximation does not always encompass the extinction of the two selected non-spherical particles in Fig. 13a and c.

## 4.2 Passive Microwave Response

In addition to the radar sensitivity to melting, there is also an interest in understanding the sensitivity of passive microwave brightness temperatures to the onset of melting. In reality, the melting layer of a precipitating cloud is likely to be obscured by an overlying ice region, obscuring it partially or wholly, depending on the wavelength of radiation. For the following simplified analysis, a single layer of melting hydrometeors is simulated, with no atmospheric gases or other layers intervening. To compute brightness temperatures, a two-stream approximation is used, and in past studies, (Johnson et al., 2012), has been found to be remarkably accurate at describing thermal emission from a plane-parallel slab viewed at the  $55^\circ$  incidence angle, typical of a satellite microwave imager. The details of the two-

stream approximation are provided in (Petty, 2008). Under the assumption of an infinitely thick slab, the transmittance through the slab is zero. Consequently the resulting upwelling brightness temperature from the two-stream approximation is only a function of the layer temperature  $T_a$ , and the layer reflectivity,  $r_\infty$ .

$$T_{B,\infty} = T_a - T_a r_\infty, \quad (5)$$

where,

$$r_\infty = \frac{(\sqrt{1 - g\tilde{\omega}} - \sqrt{1 - g})}{(\sqrt{1 - g\tilde{\omega}} + \sqrt{1 - g})}. \quad (6)$$

Here  $\tilde{\omega}$  and  $g$  have been appropriately integrated over the exponential particle size distribution (Eq. 2). In this simplified model,  $r_\infty$  is uniquely determined given the following physical properties: the wavelength of incident radiation, melting fraction (i.e., the dielectric properties), the particle orientation, and the particle size distribution. The utility of the two-stream model should not be oversold: it is useful method for understanding the bulk sensitivity of upwelling microwave TBs to modifications in the underlying physical properties of hydrometeors, but ignores all other contributions that would normally be present in an actual remote sensing scene. In this sense, it provides a “worst-case” scenario for the influence of microphysical properties on the upwelling brightness temperature.

Figure 15 shows a surface plot of two-stream brightness temperatures computed at 89, 165, and 183.31 GHz. The  $y$  axis is  $D_0$ , and the  $x$  axis is the melt fraction from 0 to 0.15. Given a fixed  $D_0$  (i.e., a fixed ice water content), the onset of melting can have a rapid and dramatic impact on the simulated TBs. For example, with  $D_0$  around 0.03 cm, proceeding from an unmelted particle to a melt fraction of 0.01 (only 1 % melted), the brightness temperatures increase by almost 100 K. Whereas for larger particles (higher ice water contents), the sensitivity to melting is not as rapid but nevertheless significant throughout the range of  $D_0$  values. The onset of melting, from 0.0 to 0.15 melt fractions, leads to increases of up to approximately 125 K. In order to better understand the impact of this, further analysis is needed wherein a 1-D vertical or 3-D simulation is used to simulate the a more realistic



atmospheric column. It is expected that an ice layer, overlaying melting and rain layer would partially or wholly obscure emission through scattering at the frequencies examined here – this is left for future studies.

## 5 Conclusions

In this paper, the Single Particle Melting Model (SPMM) was introduced as a novel and computationally efficient method to simulate the melting of an arbitrarily shaped ice hydrometeor. SPMM uses a novel nearest-neighbor method for determining when a particular point will melt and when previously melted points can move. It is easy to implement and map into any existing thermodynamic/melting-layer model, where a simple mapping between melt-water generation in a thermodynamic model and the melt-fractions generated by the melting simulation are linked as appropriate. This also provides finer control over the particle size distribution within the melting layer.

A limited study of the onset of melting, for melt fractions ranging from 0 to 0.15 was performed in order to quantify the sensitivity of microwave radiation scattering and extinction. Two snowflake aggregate shapes were selected; one composed of needles and the other composed of dendritic crystals. For comparison with past studies, the scattering and extinction properties of spherical particles having 1, 10, 50, and 100 % volume fractions of ice (relative to air) were simulated. Single particle calculations using the discrete dipole approximation, via the DDSCAT code, were made highlighting the individual particle and size-distribution integrated particle properties. These calculations were made at 13.4, 35.6, and 94 GHz; consistent with the GPM dual frequency precipitation radar, and at 89, 165, and 183.31 GHz consistent with the GPM microwave radiometer. We believe that this study represents the first simulation of the scattering and extinction properties of realistically-shaped melting hydrometeors for a wide range of microwave frequencies and particle sizes.

There is a significant sensitivity of the computed extinction and scattering properties to the base hydrometeor shape and to the onset of melting. We found, in particular, that the spherical particle assumption was unable to capture the range of computed scattering prop-

erties from the non-spherical particles, and did not provide consistent relationships between scattering and extinction throughout the onset of melting. The conclusion one could draw from this is that from a modeling perspective, it appears that spherical particles (no matter how the density/mass is modified) cannot fully represent the range of uncertainties in a lack of knowledge of the hydrometeors present in a given remote sensing scene. Capturing this behavior in physical models is critical for accurately computing uncertainty estimates in forward model simulations and retrieval algorithms.

Validation of these simulations could be, in future work, performed by examining, for example, radar observations of stratiform melting layers – especially in cases where in-situ observations of particle shape are available. The present model is currently being adapted to simulate melting for existing ice-hydrometeor databases consisting of tens of thousands of particles, and will allow for more realistic comparisons with observational data.

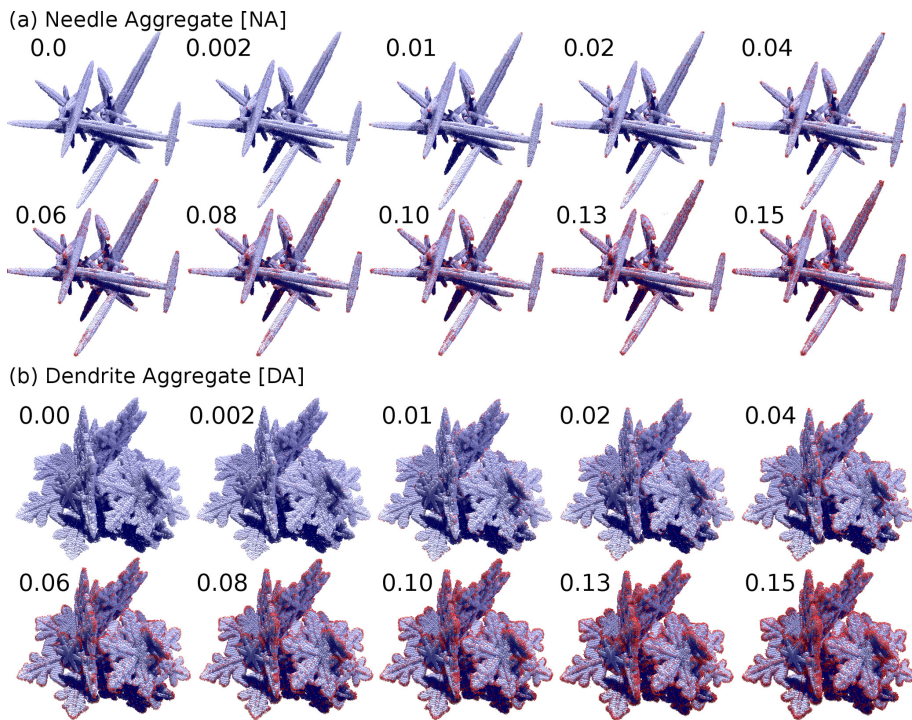
*Acknowledgements.* The majority of this investigation was supported by the NASA PMM (R. Kakar) and RST (L. Tsaoussi) programs. Specifically, NASA grants NNX11AR55G (PI: B. Johnson), NNX10AI49G (PI: W. Olson), and NNX10AT36A (PI: G. Skofronick-Jackson)

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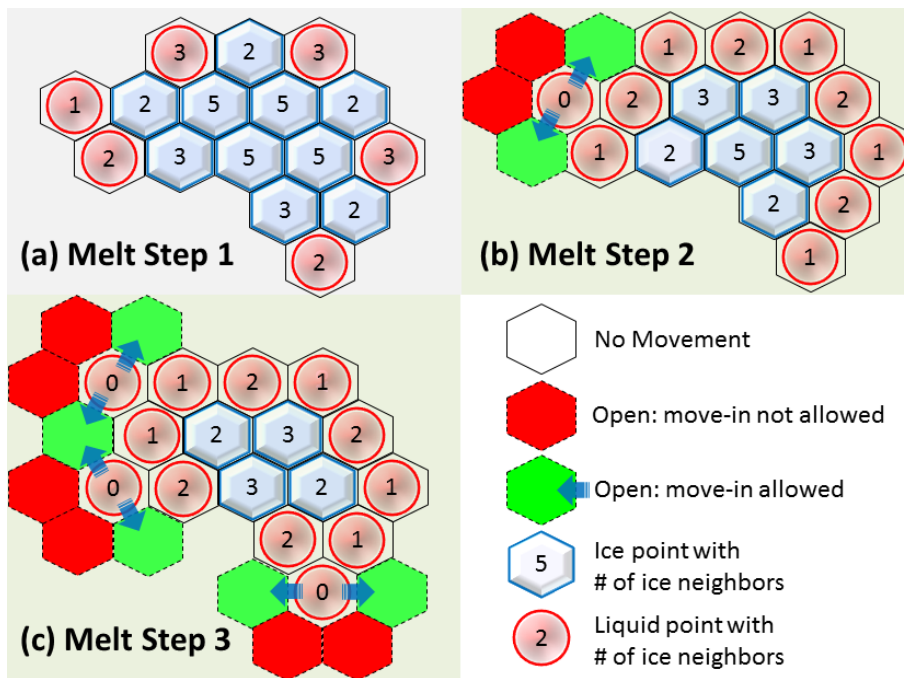
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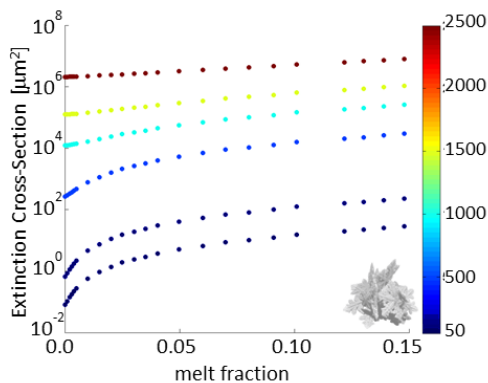


**Figure 1.** Selected steps in the onset of melting for **(a)** needle-aggregate (NA), and **(b)** dendrite-aggregate (DA). Blue regions are ice, red regions are liquid water; mass fraction of melt water is indicated at each step. Particle mass is conserved throughout the melting process.

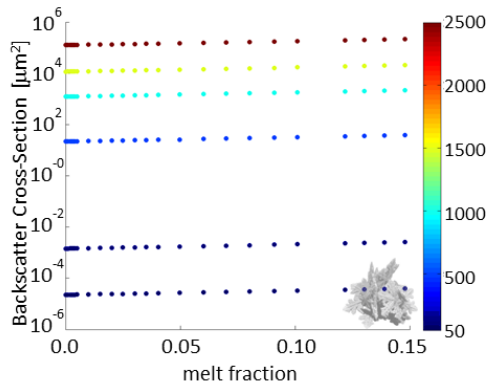


**Figure 2.** A simplified diagram depicting the steps in the SPMM. At each step, ice (blue) points are converted to liquid (red) following nearest neighbor rules (see text). Numbers depict the number of nearest ice neighbors. Panel **(a)** indicates the first iteration of the melting algorithm, **(b)** shows the second iteration, and the allowable move conditions, and **(c)** shows the third iteration, where more ice has melted and three points are now free to move into the green regions.

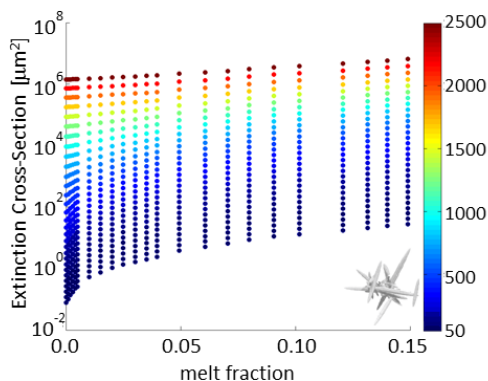
(a) [DA] 13.4 GHz Extinction Cross-Section



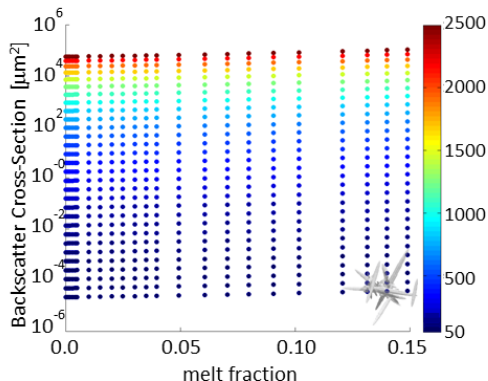
(b) [DA] 13.4 GHz Backscattering Cross-Section



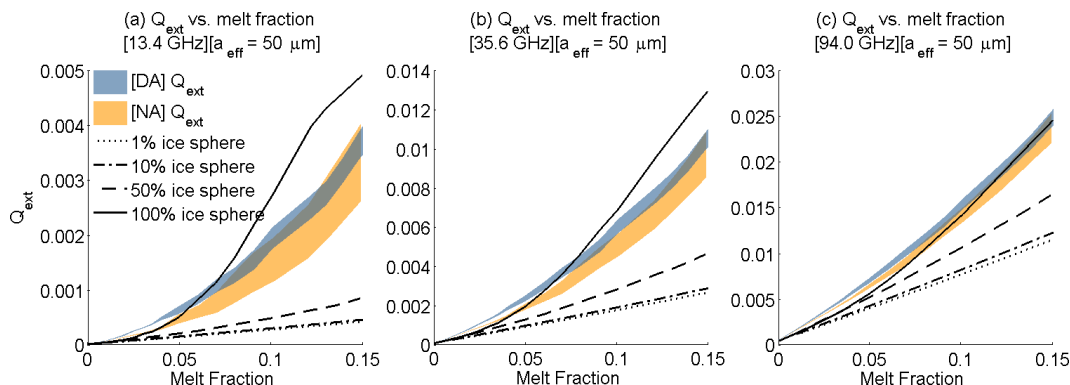
(c) [NA] 13.4 GHz Extinction Cross-Section



(d) [NA] 13.4 GHz Backscattering Cross-Section

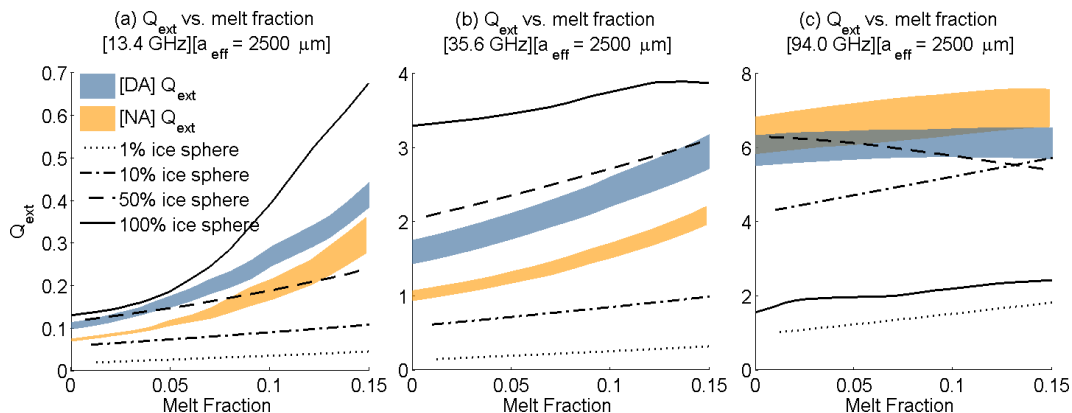


**Figure 3.** Overview of 13.4 GHz extinction and backscattering: dendrite-aggregate (row 1) (a) mean extinction cross-section vs. melt fraction and (b) mean backscattering cross-section vs. melt fraction for various values of effective radius in microns (colorbar). Needle-aggregate (row 2) (c) mean extinction cross-section vs. melt fraction, and (d) mean backscattering cross-section.

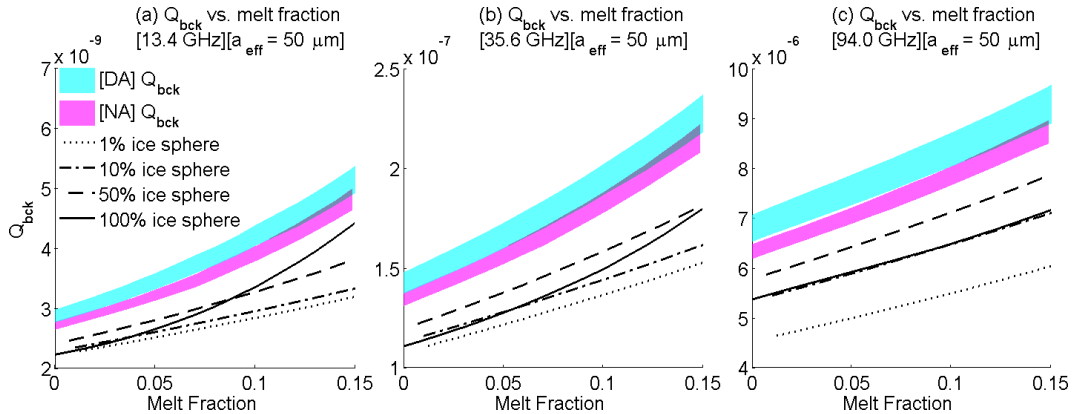


**Figure 4.** Single-particle DDSCAT calculations of the extinction coefficient ( $Q_{\text{ext}}$ ) for the dendrite-aggregate (DA) and needle-aggregate (NA) at 50 micron effective radius. Shaded regions represent the range of  $Q_{\text{ext}}$  due to various orientations of the particle relative to the direction of incident radiation. Black lines are the equivalent ice sphere  $Q_{\text{ext}}$  values at 1, 10, 50, and 100 % densities (100 % is the bulk density of ice,  $917 \text{ kg m}^{-3}$ ). **(a)** 13.4, **(b)** 35.6, **(c)** 94.0 GHz – note different vertical axis ranges on each panel for optimal visualization.

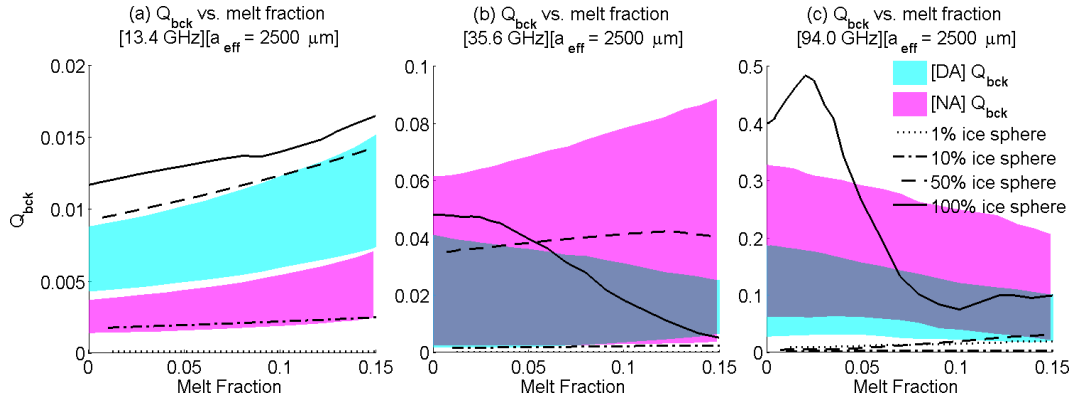




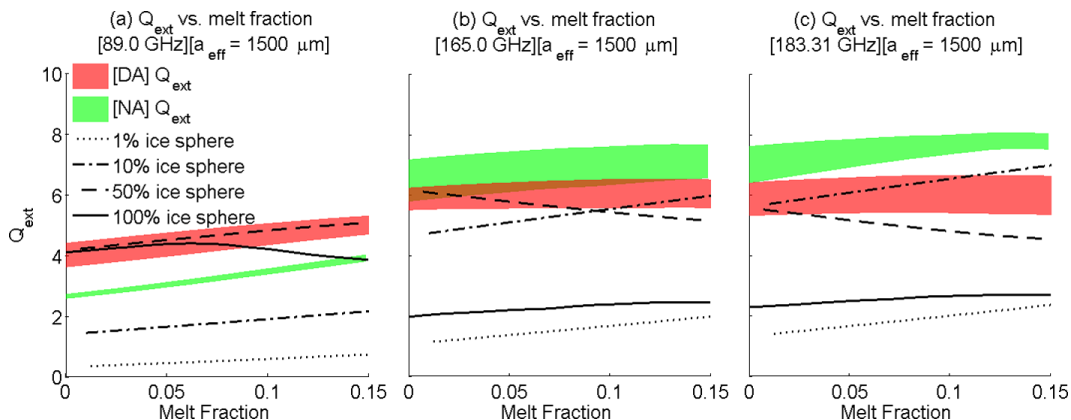
**Figure 5.** Same as Fig. 4, except effective radius is now 2500 microns (50 times larger radius; 125 000 times larger mass). As before, the vertical axes are on different scales. Notice the swap at 94 GHz, where the needle-aggregate exhibits a higher extinction than the dendrite-aggregate.



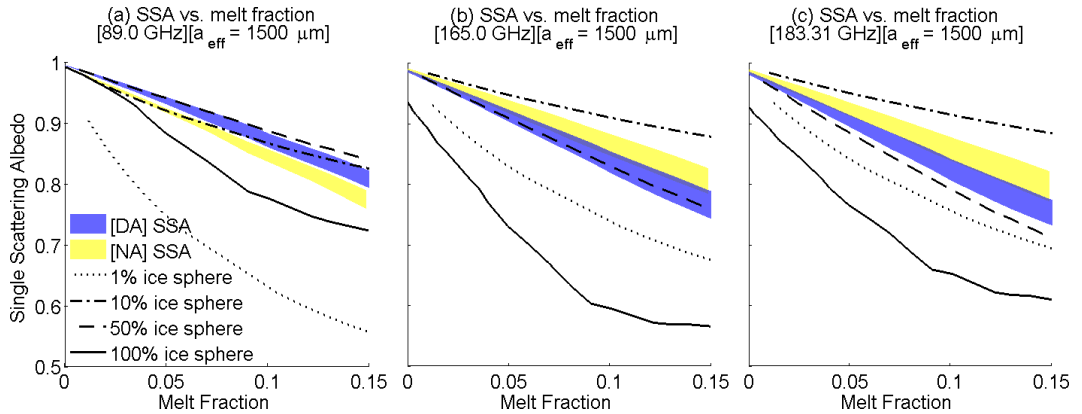
**Figure 6.** Similar to Fig. 4, the single-particle backscattering efficiency coefficient ( $Q_{bck}$ ) is shown for 50 micron effective radius. Both aggregates show consistently larger backscattering efficiencies for the same mass as melt fractions compared to the spheres. Note the difference in the vertical axes on each panel.



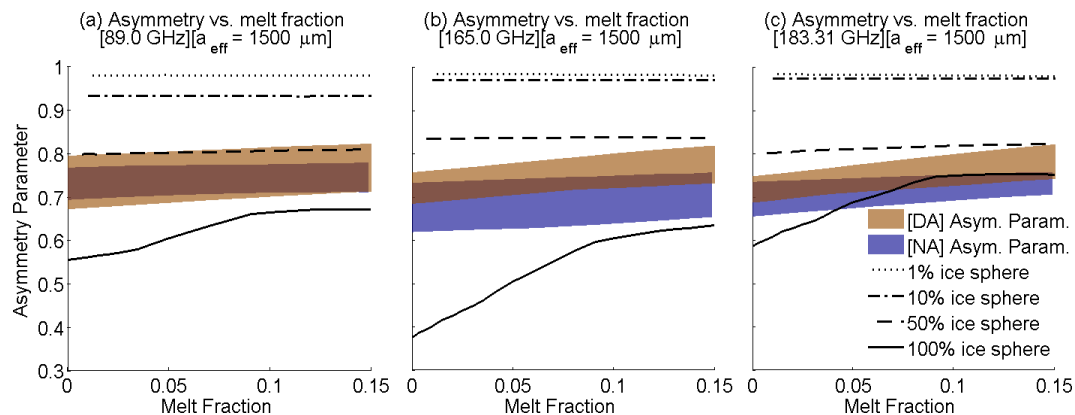
**Figure 7.** Same as Fig. 6, except at 2500 micron effective radius. The needle-aggregate at 35 GHz has a very large range due to particle rotation, whereas the dendrite-aggregate exhibits about half the same range. Similar to Fig. 6c, the needle-aggregate exhibits a larger backscattering compared to the dendrite-aggregate efficiencies at 35.6 and 94 GHz. Panel (a) shows behavior consistent with spheres, however, in (b) 35.6, and (c) 94 GHz there is no consistent behavior. Note that the backscattering efficiency in panel (b) for (DA) (blue region) is completely encompassed by (NA), and there is significant overlap between the two in panel (c).



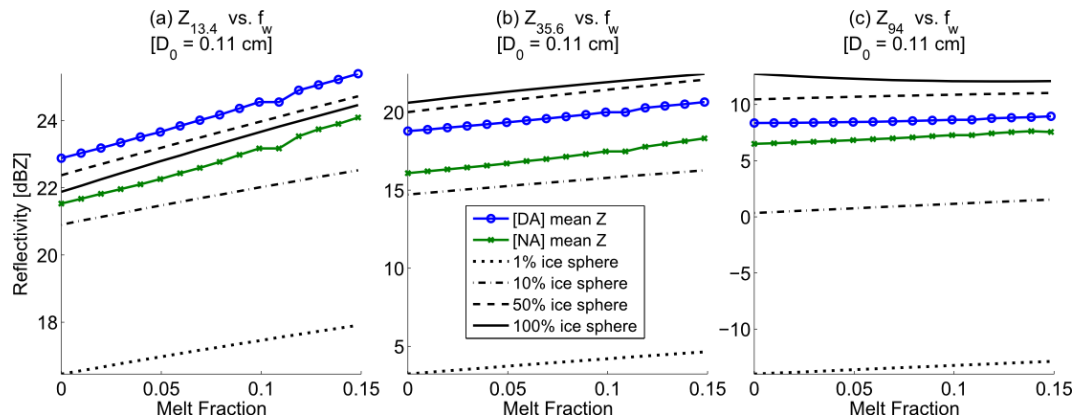
**Figure 8.** Extinction efficiencies ( $Q_{\text{ext}}$ ) computed for both aggregates and spheres at an effective radius of 1500 microns for (a) 89, (b) 165, and (c) 183.31 GHz. Similar to previous plots, shaded regions represent the ranges of  $Q_{\text{ext}}$  due various orientations of the particle relative to incident radiation. Vertical axes are on the same scale.



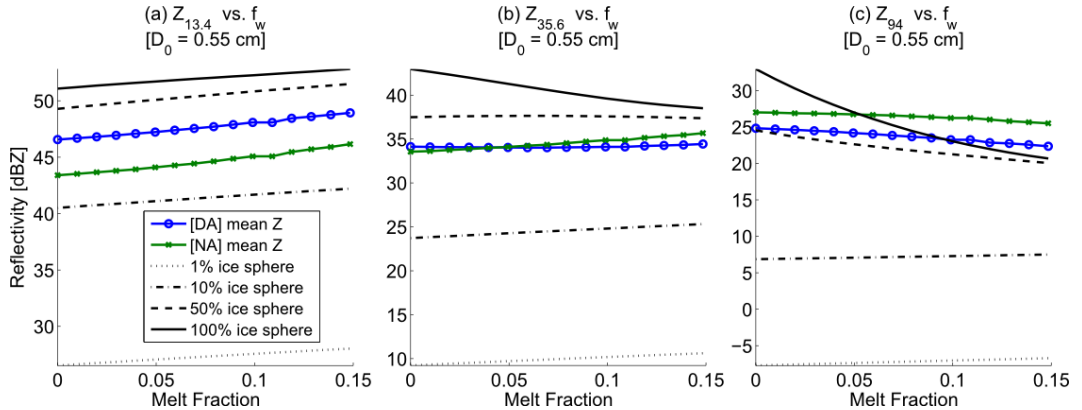
**Figure 9.** Similar to Fig. 8, except plotting single scattering albedo ( $\tilde{\omega} = Q_{\text{sca}}/Q_{\text{ext}}$ ). See text for discussion.



**Figure 10.** Following Figs. 8 and 9, the scattering asymmetry parameter (indicating the degree of forward scattering) is shown here. At all frequencies, the asymmetry parameter is relatively insensitive to the early stages of melting, consistent with the fact that the overall structure of the melting particles does not change too much over this range. In panel (a), the asymmetry for (NA) (blue region) is fully encompassed by the asymmetry from (DA), with partial overlaps in (b, c).

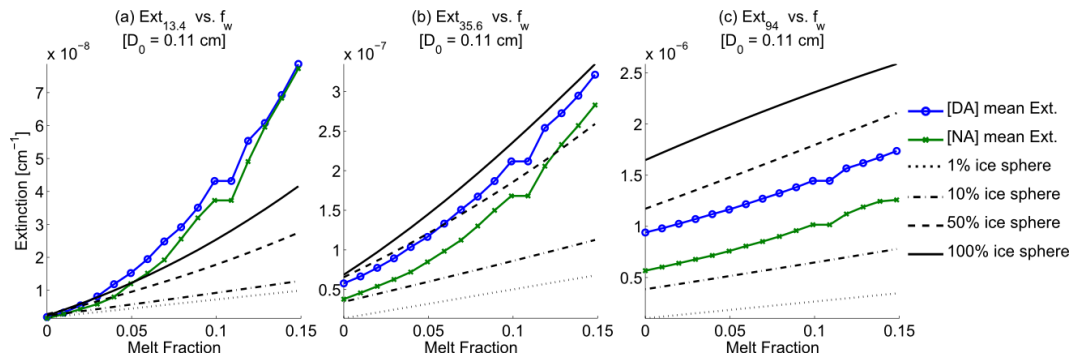


**Figure 11.** Simulated radar reflectivities at **(a)** 13.4, **(b)** 35.6, and **(c)** 94 GHz. Reflectivities were computed by integrating over the orientation-averaged single-particle backscattering cross-sections, assuming an exponential particle size distribution, for both aggregates and spheres. The reflectivities are presented only for one value of  $D_0 = 0.11$  cm (i.e., the liquid-equivalent median volume diameter).  $N_0$  and  $D_0$  are computed following Sekhon and Srivastava, 1970 (see text). Note different scaling on the vertical axes.

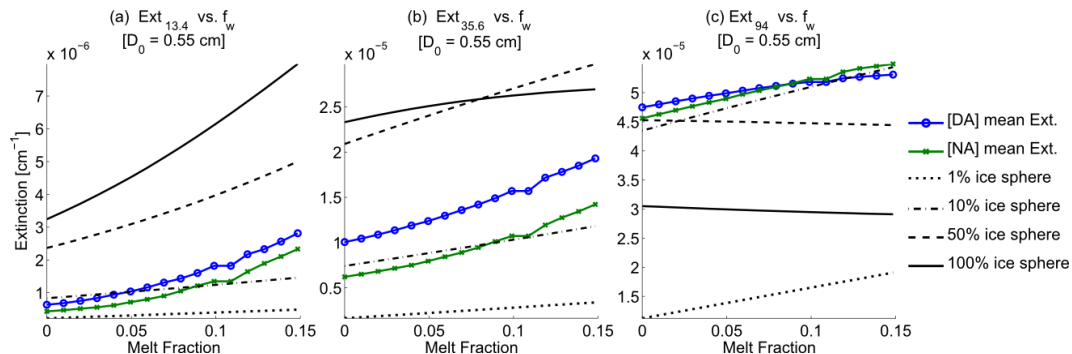


**Figure 12.** Similar to Fig. 11, except for a larger  $D_0 = 0.55$  cm. Notice the flip in the computed reflectivities at 94 GHz, similar to what was observed in the single-particle backscattering properties (Fig. 7). Note: The results in this figure may be subject to large particle diameter truncation, as discussed in the text.

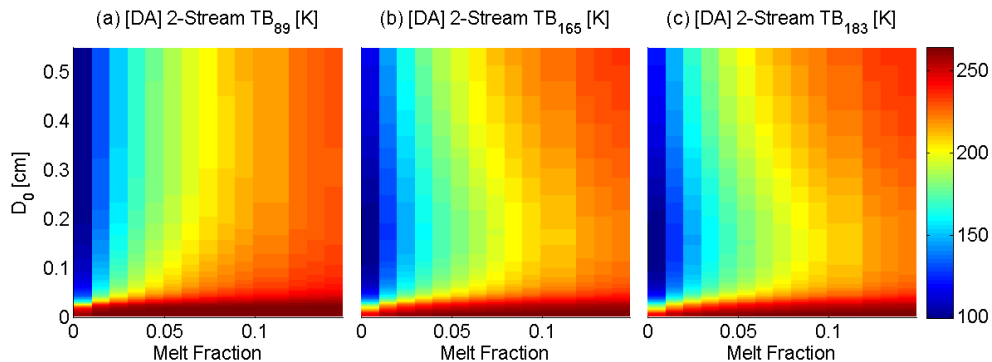




**Figure 13.** Volume extinction coefficient for  $D_0 = 0.11 \text{ cm}$  at (a) 13.4, (b) 35.6, and (c) 94.0 GHz. The orientation-averaged single-particle extinction cross-section was integrated over the same particle size distribution as was used in Fig. 11. Note different scaling on the vertical axes, including the exponents.



**Figure 14.** Same as Fig. 13 except for  $D_0 = 0.55$  cm. Note the different scaling on the vertical axes, including the exponents; and the results in this figure may be subject to large particle diameter truncation, as discussed in the text.



**Figure 15.** A surface plot of simulated brightness temperatures (TBs) for  $D_0$  vs. melt fraction assuming an “infinite”-layer of melting hydrometeors at **(a)** 89, **(b)** 165.0, and **(c)** 183.31 GHz. A two-stream radiative transfer model (see text) was used to compute the TBs. Particle size distributions used here are the same as those used in the reflectivity computations.